

Empirical Atomic Level Simulations for QC Applications

Paul von Allmen

*Jet Propulsion Laboratory
California Institute of Technology*

Contributors:

Seungwon Lee, Fabiano Oyafuso, Gerhard Klimeck

Jet Propulsion Laboratory

Tim Boykin

Univ. of Alabama, Huntsville

Collaborators:

Bob Joynt, Mark Friesen

Univ. of Wisconsin, Madison

Charter

- 1) Build simulation tools that help in designing QC in solid state implementation
- 2) Assist experimental efforts in characterizing and designing QC structures

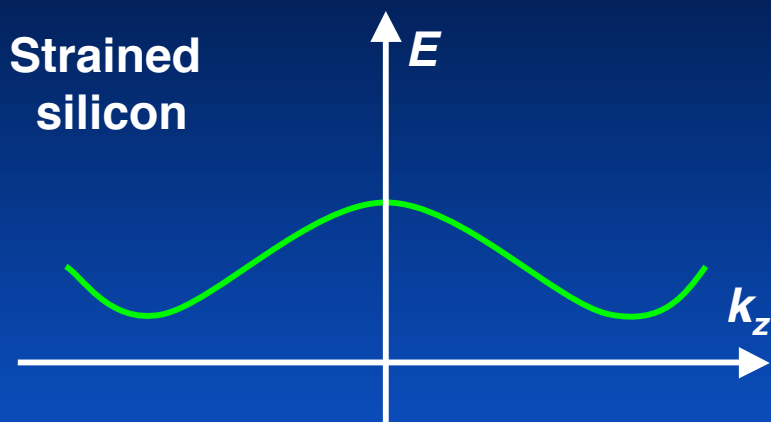
Tools

- 1) Empirical tight binding electronic structure calculation software that can simulate systems with millions of atoms, strain, Coulomb interaction, magnetic field, optical properties
NEMO-3D
- 2) Quantum transport simulator using non-equilibrium Green function formalism; phonon scattering, interface roughness, alloy disorder, atomic level resolution, spin effects (under construction) and magnetic field (under construction)
NEMO-1D and extension

Outline

1. Atomic level simulation of SiGe
2. Quantum confinement induced intervalley coupling in Si
3. Electrostatically defined quantum dot - 3D atomic level
4. Spin transport with non-equilibrium Green function
5. Atomic level simulation of quantum dot in magnetic field

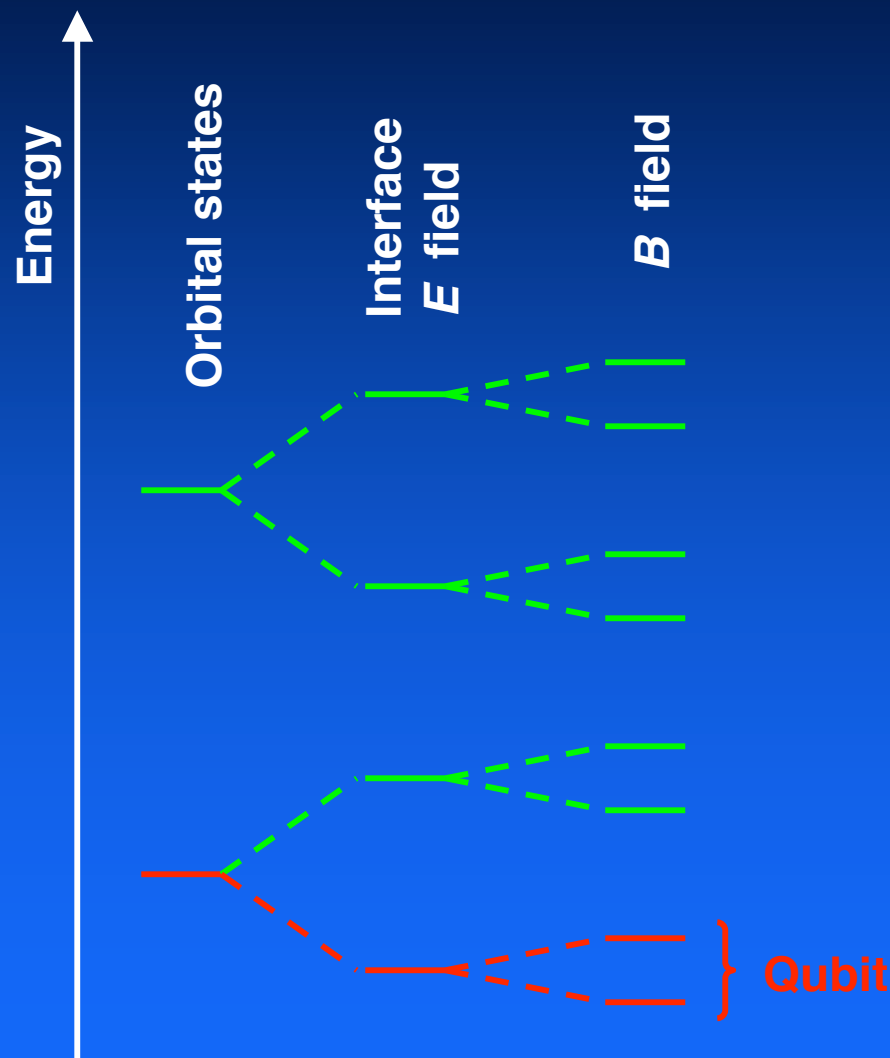
Quantum confinement induced intervalley coupling in Si



Valley degeneracy in X direction is broken by interface and electric field

Are the qubit states separated enough from higher energy states?

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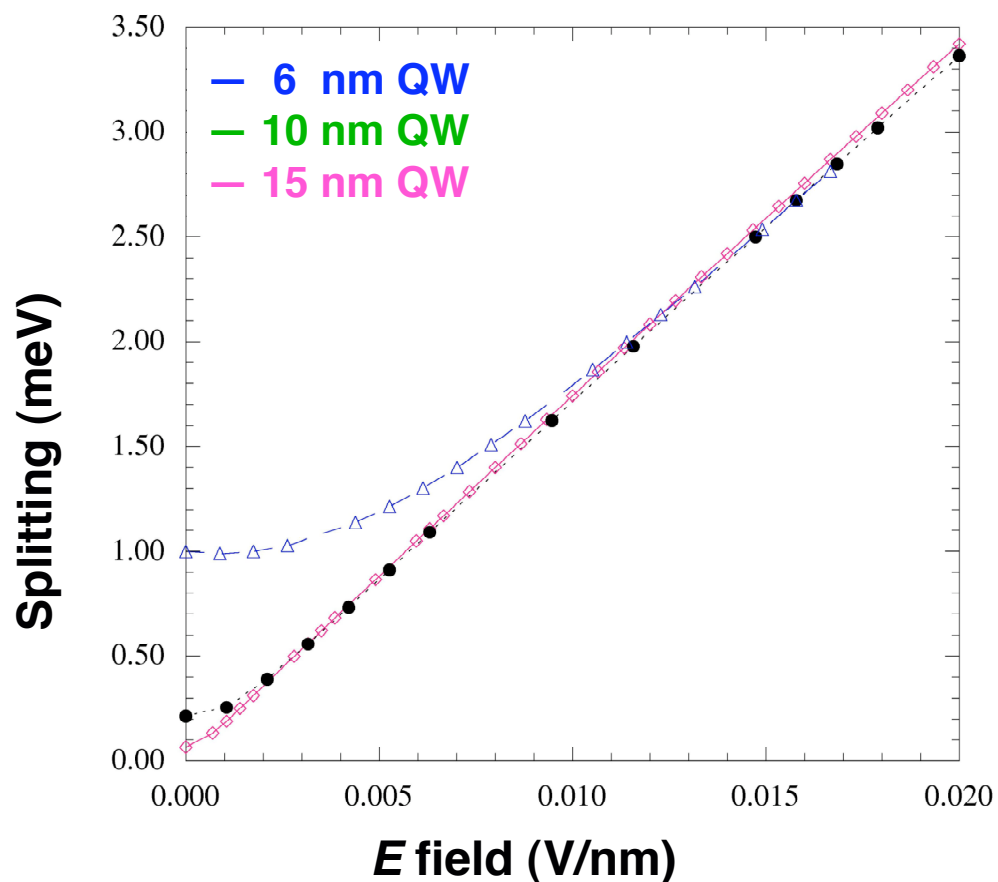


Quantum confinement induced intervalley coupling in Si

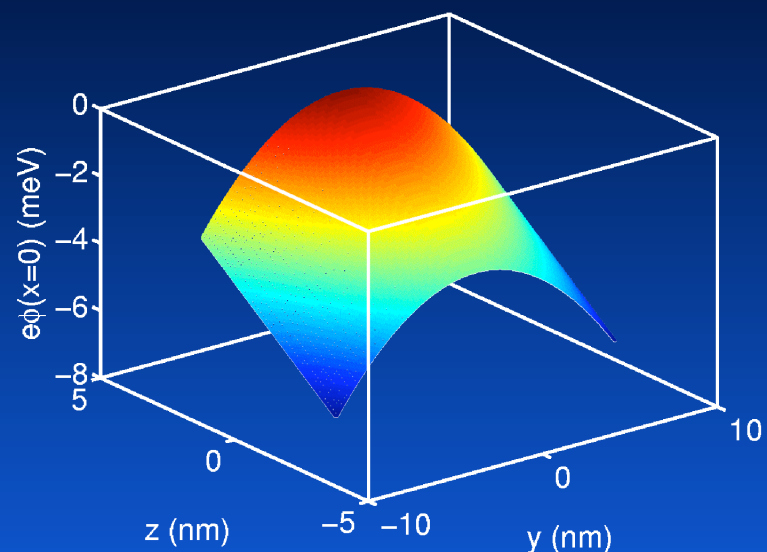
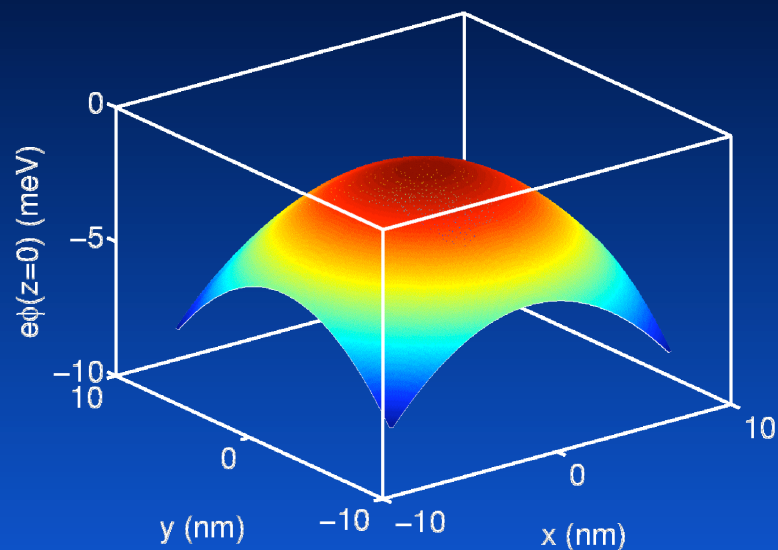
- 1D Tight Binding (NEMO 1-D)
- Si quantum well
- Strained on $\text{Si}_{0.8}\text{Ge}_{0.2}$
- Hard walls

Splitting is large enough (0.5-1.5 meV) for quantum computing

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Electrostatic confinement in Si QW

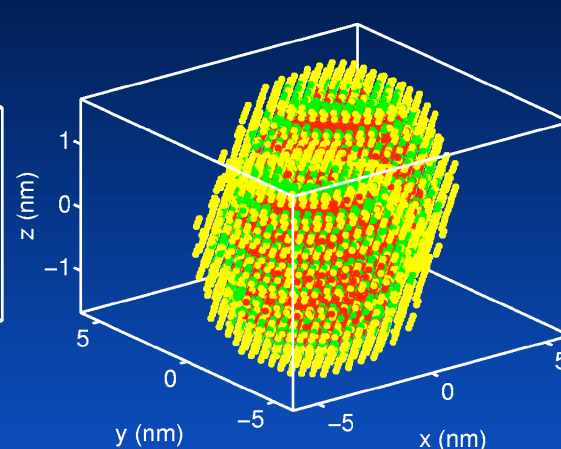
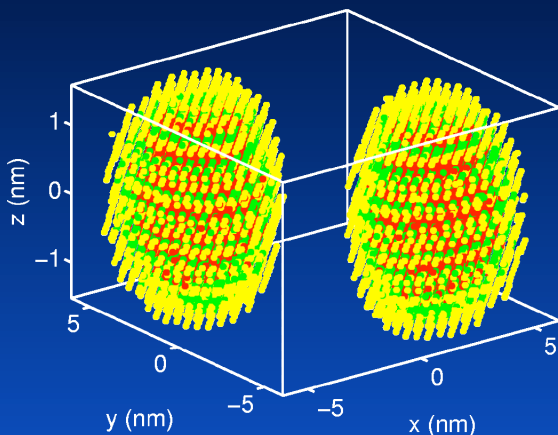
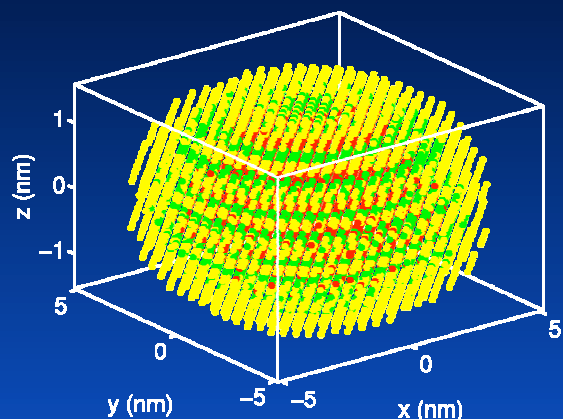


Geometry

- 5nm Si quantum well lattice-matched to $\text{Si}_{0.8}\text{Ge}_{0.2}$ (0.8% tensile strain)
- Non-zero field along z and lateral electrostatic confinement.
- Rescaled problem size (5×10^4 atoms) from WI group

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Electronic structure of low-lying states



Splitting of electronic levels

- Valley splitting due to loss of translational symmetry in QD (~ 0.45 meV)
- Zeeman splitting is smaller (~ 0.12 meV).

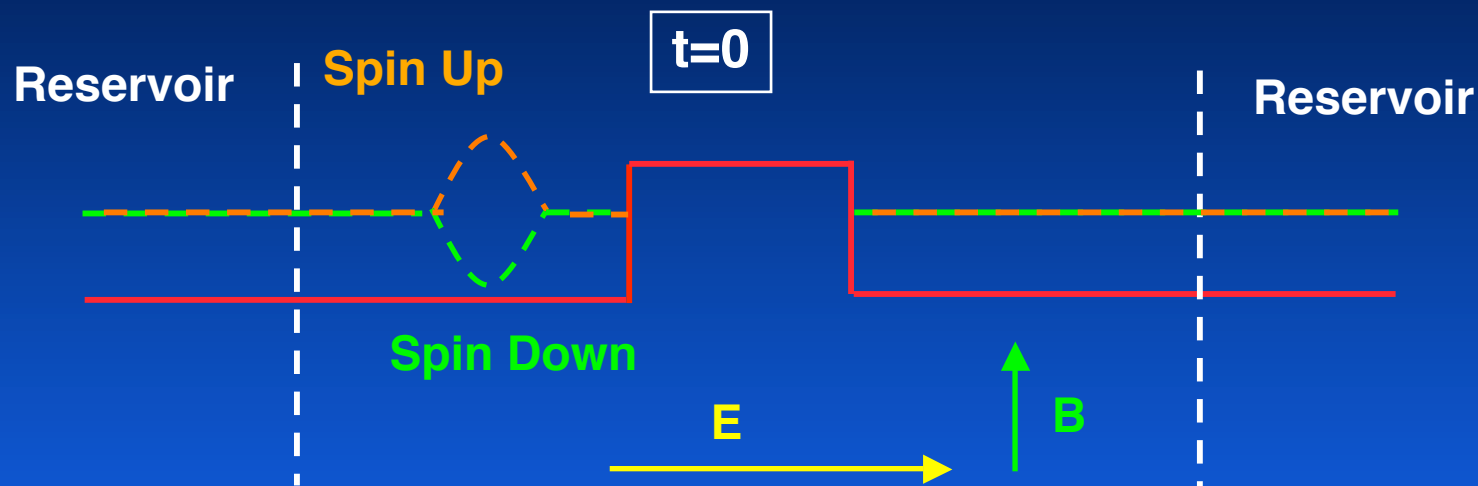
QD symmetry breaking

$B = 1$ T



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Spin Transport with Non-Equilibrium Green Functions



Time evolution

1. Spin transport
2. Spin precession
3. Spin relaxation and decoherence

Applications

1. Magneto-optic (TRKR) experiments
2. Transport between qubits
3. Preparation of initial state for QC

Spin Transport with Non-Equilibrium Green Functions

Equations of motion

$$\begin{aligned}
 (i\partial_t - H^{(0)})G^<(xt, x|t) &= \int dx' \int dt' \delta(x - x') \delta(t - t') G^a(x'|t', x|t) + \int dx' \int dt' \delta(x - x') \delta(t - t') G^<(x'|t', x|t) \\
 (i\partial_t - H^{(0)})G^r(xt, x|t) &= \delta(t - t') \delta(x - x') + \int dx' \int dt' \delta(x - x') \delta(t - t') G^r(x'|t', x|t) \\
 H^{(0)} &= \frac{p^2}{2m} + V(x) + \mathbf{B} \cdot \mathbf{S}
 \end{aligned}$$

Open boundaries (finite difference or tight binding)

$$\begin{aligned}
 (i\partial_t - H_{nm}^{(0)})G_{mn}^< &= \delta_{n1} \delta_{11}^{Br} G_{1n}^< + \delta_{n1} \delta_{11}^{B<} G_{1n}^a + \delta_{nN} \delta_{NN}^{Br} G_{Nn}^< + \delta_{nN} \delta_{NN}^{B<} G_{Nn}^a \\
 (i\partial_t - H_{nm}^{(0)})G_{mn}^r &= \delta_{nn} + \delta_{n1} \delta_{11}^{Br} G_{1n}^r + \delta_{nN} \delta_{NN}^{Br} G_{Nn}^r \\
 \delta_{11} &= H_{10}^{(0)} G_{00}^0 H_{01}^{(0)} \\
 \delta_{NN} &= H_{NN+1}^{(0)} G_{N+1N+1}^0 H_{N+1N}^{(0)}
 \end{aligned}$$

Surface Green Functions

1. Stationary in time (NEMO-1D)
 - Fourier transform of Green function
 - Algebraic Dyson equation
 - Recursive method to calculate the surface Green function
2. Time dependent case

$$G_{00}^{<0}(t, t') = g_{00}^{<0}(t - t') M^<(t, t')$$

$g_{00}^{<0}(t - t')$ is the surface Green function for a spinless system

$$M^<(t, t') = \begin{bmatrix} \cos \frac{\omega_L t}{2} \cos \frac{\omega_L t'}{2} & -i \cos \frac{\omega_L t}{2} \sin \frac{\omega_L t'}{2} \\ i \sin \frac{\omega_L t}{2} \cos \frac{\omega_L t'}{2} & \sin \frac{\omega_L t}{2} \sin \frac{\omega_L t'}{2} \end{bmatrix} \quad \text{Polarized spin up}$$

$$\omega_L = \omega_B$$

Tight-binding model for magneto-optical response

- Gauge-invariant modification of tight-binding Hamiltonian**

[M. Graf and P. Vogl, Phys. Rev. B **51**, 4940 (1995)]

$$\square_i = \square_i^0 + \square_B B \cdot \square_i$$

On-site interaction

$$t_{ij} = t_{ij}^0 \exp\left(\square \frac{ie}{\hbar c} \int_{R_j}^{R_i} A \cdot dr\right)$$

Nearest-neighbor interaction

- Absorption rate between conduction and valence levels**

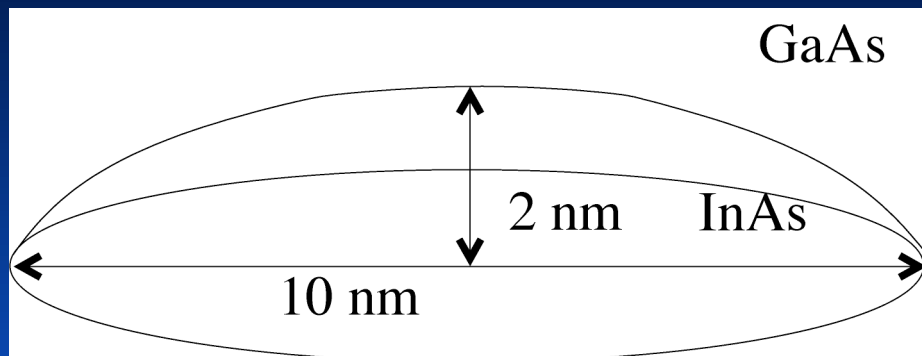
[S. Lee et al., Phys. Rev. B **66**, 235307 (2002)]

$$\square(E) = \frac{2\square}{\hbar} |\langle c | \hat{r} | v \rangle|^2 \square(E_c - E_v - E)$$

$$\langle c | \hat{r} | v \rangle = \sum_{ij} c_j^* v_i [\hat{r}_i \square_{ij} + \langle j | \square \hat{r}_i | i \rangle]$$

InAs self-assembled lens-shaped dot

- **Modeled system geometry**



- **Strain profile model: atomic elasticity model**

[P. N. Keating, Phys. Rev. B **145**, 637 (1966)]

Strain energy as a function of bond length and bond angle.

- **Tight-binding Hamiltonian**

Basis orbitals: $sp^3d^5 s^*$

Parameters generated by a genetic algorithm.

Size : 921600 = 46080 atoms X 10 basis orbitals X 2 spins

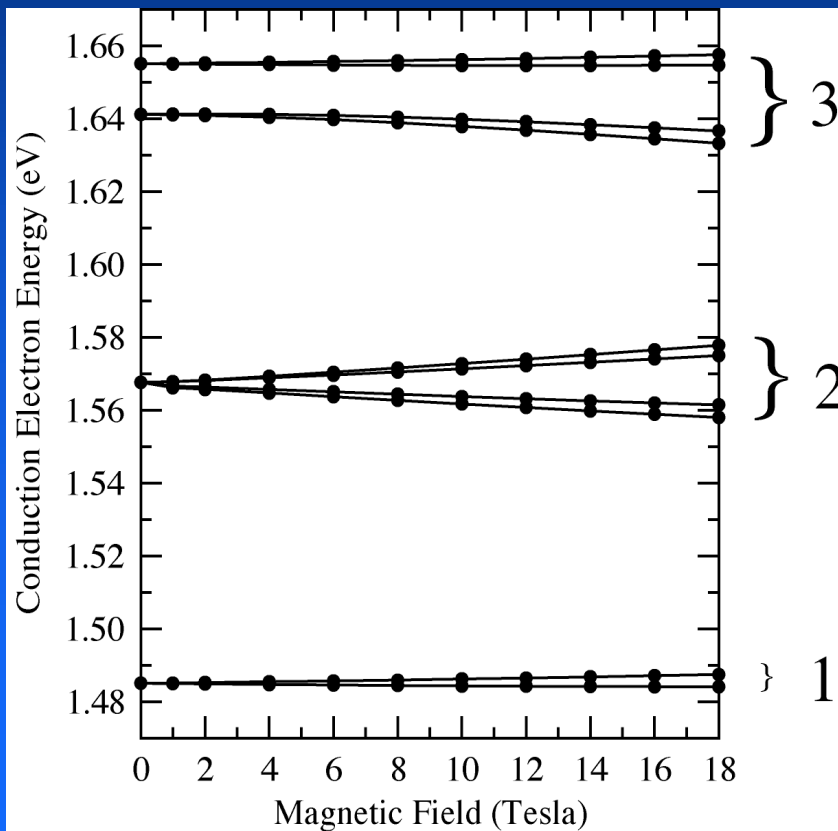
- **Eigenvalue solver: Arnoldi method with PARPACK package**

Computation on a Beowulf cluster with 30 nodes.

Magnetic-Field Effect on Electronic Structure of InAs Dot

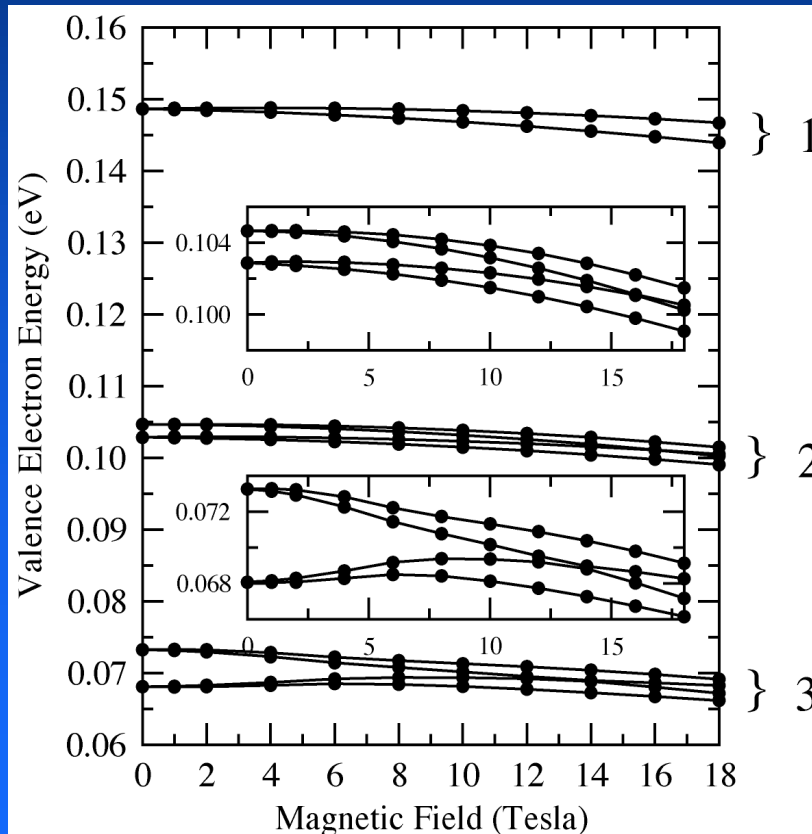
Conduction Electron Levels

- Zeeman Interaction splits the levels into spin-up and -down levels.
- Effective g-factor ($g = (E_{\uparrow} - E_{\downarrow}) / \mu_B B$) ranges from 2 to 3.5.



Valence Electron Levels

- Zeeman interaction splits the levels into $J_z = 3/2$ and $J_z = -3/2$ levels.
- Zeeman interaction couples closely-spaced levels.



Magneto-Optical Response of InAs Quantum Dot

- Selective dipole coupling between electron and hole levels.
- The selectivity remains intact even at a high magnetic field.

$$\Delta L = 0 \text{ \& \; } \Delta j = 1$$

